

### Problem 1. Edge states in the SSH chain

In this problem we will explicitly consider both an semi-infinite chain as well as a finite chain to understand both edge states and their splitting

For the semi-infinite chain, the SSH Hamiltonian takes the form

$$H = \sum_{n=1}^{\infty} (t_1 |n, A\rangle \langle n, B| + t_2 |n+1, A\rangle \langle n, B| + \text{h.c.}). \quad (1)$$

- (a) Let  $|\psi\rangle = \sum_n (\psi_{n,A} |n, A\rangle + \psi_{n,B} |n, B\rangle)$  and evaluate  $H |\psi\rangle = E |\psi\rangle$  to obtain a recurrence relation for  $\psi_{n,A}$  and  $\psi_{n,B}$ .
- (b) Derive a (normalizeable) solution for  $E = 0$ . Plot  $|\psi_{n,A}|^2$  and  $|\psi_{n,B}|^2$  vs.  $n$ . What conditions are needed for it to be normalizeable?
- (c) Write the recurrence relations in such a way that you eliminate  $\psi_{nA}$  (to obtain a recurrence relation solely in  $\psi_{nB}$ ), and using the boundary condition  $\psi_{0B} = 0$ , solve the recurrence relations (Recall that if  $a_{n+2} = pa_{n+1} + qa_n$  that we can find two solutions with  $a_n = c\lambda_1^n + b\lambda_2^n$ ). We require normalization so  $|\psi_{nB}| \rightarrow 0$  as  $n \rightarrow \infty$ , this will restrict  $|\lambda_j| \leq 1$  — in fact you should find that for  $E > 0$ ,  $|\lambda_j| = 1$ . If we let  $E = 0$ , can we recover the solution from part (b)? Why or why not?
- (d) Lastly, consider now the finite chain

$$H = \sum_{n=1}^L (t_1 |n, A\rangle \langle n, B| + t_2 |n+1, A\rangle \langle n, B| + \text{h.c.}). \quad (2)$$

Argue or compute what the edge state localized at  $L$  looks like based on what you found in part (b). Approximate the  $n = 1$  edge state with  $|\text{Left}\rangle = \sum_{n=1}^L (\psi_{nA} |n, A\rangle + \psi_{nB} |nB\rangle)$  ( $\psi_{nA}$  and  $\psi_{nB}$  coming from your part (b) solution) and similarly for the edge state localized at  $L$ , which we call  $|\text{Right}\rangle$ . With these edge states, compute the effective Hamiltonian

$$H_{\text{edge}} = \begin{pmatrix} \langle \text{Left} | H | \text{Left} \rangle & \langle \text{Left} | H | \text{Right} \rangle \\ \langle \text{Right} | H | \text{Left} \rangle & \langle \text{Right} | H | \text{Right} \rangle \end{pmatrix}. \quad (3)$$

What are the eigenstates and what is the effective gap between them?

## Problem 2. Modified Bulk SSH model

Consider the usual SSH Hamiltonian

$$H_0 = \sum_{n=-\infty}^{\infty} (t_1 |n, A\rangle \langle n, B| + t_2 |n+1, A\rangle \langle n, B| + \text{h.c.}), \quad (4)$$

but with an added term

$$V = t' \sum_{n=-\infty}^{\infty} (i |n+1, A\rangle \langle n, A| - i |n, A\rangle \langle n+1, A|) \quad (5)$$

- For the full Hamiltonian  $H = H_0 + V$  what symmetries remain in this Hamiltonian (list the symmetries of  $H_0$  and indicate which ones  $V$  breaks and which ones it preserves)? What is its topological classification ( $0$ ,  $\mathbb{Z}_2$ , or  $\mathbb{Z}$ )?
- Find the  $k$ -space Hamiltonian for  $H = H_0 + V$  and cast it into the form  $H = \epsilon(k) + \mathbf{d}(k) \cdot \boldsymbol{\sigma}$ . What are its eigenenergies?
- Compute the polarization of the bottom band.

### Problem 3. The quantum Hall effect with spin

Electrons have spin and that directly interacts with a magnetic field. In this problem, we will explore the implications of that.

- (a) First, assuming a magnetic field in the  $\mathbf{B} = B\hat{\mathbf{z}}$  in a two-dimensional electron gas. Using your favorite gauge, find the eigenstates and eigenenergies of the Pauli Hamiltonian

$$H_P = \frac{1}{2m} \left( \mathbf{p} - \frac{e}{c} \mathbf{A} \right)^2 - \frac{e}{2mc} \mathbf{B} \cdot \boldsymbol{\sigma}. \quad (6)$$

In particular, note how spin changes things: What is the degeneracy of each Landau level (make your space finite in whichever way is convenient for your gauge so that you have a total flux  $\Phi$  through the system)? How has the inclusion of spin changed this?

- (b) Define an operator  $Q_1 = \frac{1}{\sqrt{4m}} (\mathbf{p} - \frac{e}{c} \mathbf{A}) \cdot \boldsymbol{\sigma}$ , and show that  $H_P = 2Q_1^2$ . and show that if you have an eigenstate  $|\psi_E\rangle$  then  $|\psi'_E\rangle = \sqrt{2/E} Q_1 |\psi_E\rangle$  is also an eigenstate (in general or for this specific problem).  $Q_1$  is known as a *supercharge*.
- (c) *Complex supercharge*. Let the vector potential be  $\mathbf{A} = (A_x(\mathbf{r}), A_y(\mathbf{r}), 0)$  (with  $\mathbf{A}$  independent of  $z$ ; don't specify the gauge any further), and define

$$A = \frac{1}{\sqrt{2m}} \left[ \left( p_x - \frac{e}{c} A_x \right) - i \left( p_y - \frac{e}{c} A_y \right) \right], \quad Q = (\sigma_x + i\sigma_y)A. \quad (7)$$

Compute  $Q^2$ ,  $\{Q, Q^\dagger\}$ , and  $[A, A^\dagger]$ . (Recall that the cyclotron frequency is  $\omega_c = \frac{eB}{mc}$ .)

- (d) Write the Hamiltonian in the basis of spin  $\uparrow$  and spin  $\downarrow$ , and purely in terms of  $A$  and  $A^\dagger$ . Specifically, show that the partially projected operators  $\langle \uparrow | H_P | \uparrow \rangle = A^\dagger A$  and  $\langle \downarrow | H_P | \downarrow \rangle = AA^\dagger$  in terms of  $A$  and  $A^\dagger$ . These expressions are said to be *isospectral* in that the spectrum of both are the same except at zero energy. In particular, something called an *index theorem* relates them:

$$(\# \text{ of zeros of } A^\dagger A) - (\# \text{ of zeros of } AA^\dagger) = \Delta \quad (8)$$

What is the index  $\Delta$  in this problem? (*Hint*: Don't forget the degeneracy from part (a))

- (e) Given a chemical potential  $\mu = \epsilon_F = \frac{5}{2}\omega_c$  what is the magnetization  $M = \langle \sigma_z \rangle = \sum_{E < \mu} \langle \psi_E | \sigma_z | \psi_E \rangle$ ? Relate this back to  $\Delta$  from part (d).